

SELF-SIMILAR SOLUTIONS OF THE NON-STRICTLY HYPERBOLIC WHITHAM EQUATIONS FOR THE KDV HIERARCHY

V. U. PIERCE AND FEI-RAN TIAN

ABSTRACT. We study the Whitham equations for all the higher order KdV equations. The Whitham equations are neither strictly hyperbolic nor genuinely nonlinear. We are interested in the solution of the Whitham equations when the initial values are given by a step function.

1. INTRODUCTION

It is known that the solution of the KdV equation

$$(1.1) \quad u_t + 6uu_x + \epsilon^2 u_{xxx} = 0$$

has a weak limit as $\epsilon \rightarrow 0$ while the initial values

$$u(x, 0; \epsilon) = u_0(x)$$

are fixed.

This weak limit is described by hyperbolic equations. It satisfies the Burgers equation

$$(1.2) \quad u_t + (3u^2)_x = 0$$

until its solution develops shocks. Immediately after shock, the weak limit is governed by the Whitham equations [5, 6, 14, 15]

$$(1.3) \quad u_{it} + \lambda_i(u_1, u_2, u_3)u_{ix} = 0, \quad i = 1, 2, 3,$$

where the λ_i 's are given by formulae (2.12). Equations (1.3) form a 3×3 system of hyperbolic equations. After the breaking of the solution of (1.3), the weak limit is described by a 5×5 systems of hyperbolic equations similar to (1.3). Similarly, after the solution of the 5×5 system breaks down,

the weak limit is characterized by a 7×7 system of hyperbolic equations. In other words, for general initial data $u_0(x)$, one constructs the weak limit by patching together solutions of (1.2), (1.3), 5×5 , 7×7 , etc systems in the x - t plane.

The KdV equation (1.1) is just the first of an infinite sequence of equations. All these so-called higher order KdV equations can be cast in the Hamiltonian form

$$(1.4) \quad \frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \frac{\delta H_m}{\delta u} = 0, \quad m = 1, 2, \dots,$$

where H_m 's form a sequence of conserved functionals for the KdV equation. The small dispersive parameter ϵ is hidden in H_m . In particular, when $m = 1$, (1.4) is the KdV equation.

The solution of the higher order KdV equation (1.4) also has a weak limit as $\epsilon \rightarrow 0$. As in the KdV case, this weak limit satisfies the Burgers type equation

$$(1.5) \quad u_t + \left(\frac{C_m}{m+1} u^{m+1} \right)_x = 0,$$

where C_m is given in (2.20), until the solution of (1.5) forms a shock. After the breaking of the solution of (1.5), the limit is governed by equations similar to (1.3), namely,

$$(1.6) \quad u_{it} + \mu_i^{(m)}(u_1, u_2, u_3) u_{ix} = 0, \quad i = 1, 2, 3,$$

where $\mu_i^{(m)}$'s are given in (2.15). They will also be called the Whitham equations. As in the KdV case, after the solution of (1.6) breaks down, the weak limit is described by a 5×5 system of hyperbolic equations.

In this paper, we are interested in the solution of the Whitham equations for the higher order KdV (1.4) with a step-like initial function

$$(1.7) \quad u_0(x) = \begin{cases} 1 & x < 0 \\ 0 & x > 0. \end{cases}$$

For such an initial function, the solution of the Burgers type equation (1.5) has already developed a shock at the initial time, $t = 0$. Hence, immediately after $t = 0$, the Whitham equations (1.6) kick

in. Solutions of (1.6) occupy some domains of the space-time while solutions of (1.5) occupy other domains. These solutions are matched on the boundaries of the domains.

Equations (1.2) and equations (1.5) are prototypes in the theory of hyperbolic conservation laws [7]. Their solutions will generally develop shocks in finite times. The solutions can be extended beyond the singularities as the entropy solutions.

Solutions of equations (1.2) or equations (1.5), in the theory of the zero dispersion limit, are not extended as weak or entropy solutions after the formation of singularities. Instead, they are extended to match the Whitham solutions of (1.3) or (1.6). For initial data (1.7), the resulting solutions of the Whitham equations (1.6) will be seen to be more complex than those of (1.3) in the KdV case.

The KdV case with the step-like initial data (1.7) was first studied by Gurevich and Pitaevskii [3]. They found that it was enough to use the Burgers solution of (1.2) and the Whitham solution of (1.3) to cover the whole x - t plane, without going to the 5×5 or 7×7 system. Namely, the space-time is divided into three parts

$$(1) \frac{x}{t} < -6, \quad (2) -6 < \frac{x}{t} < 4, \quad (3) \frac{x}{t} > 4.$$

The solution of (1.2) occupies the first and third parts,

$$(1.8) \quad u(x, t) \equiv 1 \quad \text{when } \frac{x}{t} < -6, \quad u(x, t) \equiv 0 \quad \text{when } \frac{x}{t} > 4.$$

The Whitham solution of (1.3) lives in the second part,

$$(1.9) \quad u_1(x, t) \equiv 1, \quad \frac{x}{t} = \lambda_2(1, u_2, 0), \quad u_3(x, t) \equiv 0,$$

when $-6 < x/t < 4$.

Whether the second equation of (1.9) can be inverted to give u_2 as a function of the self-similarity variable x/t hinges on whether

$$\frac{\partial \lambda_2}{\partial u_2}(1, u_2, 0) \neq 0.$$

Indeed, Levermore [8] has proved the genuine nonlinearity of the Whitham equations (1.3), i.e.,

$$(1.10) \quad \frac{\partial \lambda_i}{\partial u_i}(u_1, u_2, u_3) > 0, \quad i = 1, 2, 3,$$

for $u_1 > u_2 > u_3$.

For the higher order KdV (1.4), equations (1.6), in general, are not genuinely nonlinear, i.e., a property like (1.10) is not available. Hence, solutions like (1.8) and (1.9) need to be modified.

Our construction of solutions of the Whitham equation (1.6) makes use of the non-strict hyperbolicity of the equations. For KdV, it is known that the Whitham equations (1.3) are strictly hyperbolic, namely:

$$\lambda_1(u_1, u_2, u_3) > \lambda_2(u_1, u_2, u_3) > \lambda_3(u_1, u_2, u_3)$$

for $u_1 > u_2 > u_3$ [8]. For the higher order KdV (1.4), different eigenspeeds of (1.6), $\mu_i^{(m)}(u_1, u_2, u_3)$'s, may coalesce in the region $u_1 > u_2 > u_3$ [9].

For the higher order KdV with step-like initial function (1.7), the space time is divided into four regions (see Figure 1.)

$$(1) \frac{x}{t} < \alpha, \quad (2) \alpha < \frac{x}{t} < \beta, \quad (3) \beta < \frac{x}{t} < 4^m, \quad (4) \frac{x}{t} > 4^m,$$

where α and β are some constants. In the first and fourth regions, the solution of (1.5) governs the evolution:

$$u(x, t) \equiv 1 \quad \text{where } x/t < \alpha \text{ and } u(x, t) \equiv 0 \quad \text{where } x/t > 4^m.$$

The Whitham solution of (1.6) lives in the second and third regions; namely:

$$(1.11) \quad u_1(x, t) \equiv 1, \quad \frac{x}{t} = \mu_2^{(m)}(1, u_2, u_3), \quad \frac{x}{t} = \mu_3^{(m)}(1, u_2, u_3),$$

when $\alpha < x/t < \beta$, and

$$(1.12) \quad u_1(x, t) \equiv 1, \quad \frac{x}{t} = \mu_2^{(m)}(1, u_2, 0), \quad u_3(x, t) \equiv 0,$$

when $\beta < x/t < 4^m$.

Equations (1.11) yield

$$\mu_2^{(m)}(1, u_2, u_3) = \mu_3^{(m)}(1, u_2, u_3)$$

on a curve in the region $0 < u_3 < u_2 < 1$. This implies the non-strict hyperbolicity of the Whitham equations (1.6) for the KdV hierarchy.

The $m = 2$ case has been studied in [9]. There, inequalities

$$(1.13) \quad \frac{\partial \mu_3^{(m)}}{\partial u_3} < \frac{3}{2} \frac{\mu_2^{(m)} - \mu_3^{(m)}}{u_2 - u_3} < \frac{\partial \mu_2^{(m)}}{\partial u_2} \quad \text{for } u_1 > u_2 > u_3 > 0$$

have played a crucial role in verifying that equations (1.11) or (1.12) can indeed be solved to give the solution of the Whitham equations (1.6) when $m = 2$.

For $m > 2$, inequalities (1.13) are not valid any more. We therefore must use a different approach to solve the problem. The calculations are considerably more difficult than in the $m = 2$ case. This is mainly because q of (2.25) is a polynomial of degree m when $m > 2$ while it is only a quadratic polynomial when $m = 2$.

The organization of the paper is as follows. In Section 2, we will study the eigenspeeds, $\mu_i^{(m)}$'s, of the Whitham equations (1.6). In Section 3, we will construct the self-similar solution of the Whitham equations for the initial function (1.7). In Section 4, we will use the self-similar solution of Section 3 to construct the minimizer of a variational problem for the zero dispersion limit of the KdV hierarchy.

In a subsequent publication, we will study the Whitham solutions for all the other step-like initial data.

2. THE WHITHAM EQUATIONS

In this section we define the eigenspeeds of the Whitham equations for both the KdV (1.1) and higher order KdV (1.4). We first introduce the polynomials of ξ for $n = 0, 1, 2, \dots$ [1, 4, 12]:

$$(2.1) \quad P_n(\xi, u_1, u_2, u_3) = \xi^{n+1} + a_{n,1}\xi^n + \dots + a_{n,n+1} ,$$

where the coefficients, $a_{n,1}, a_{n,2}, \dots, a_{n,n+1}$ are uniquely determined by the two conditions

$$(2.2) \quad \frac{P_n(\xi, u_1, u_2, u_3)}{\sqrt{(\xi - u_1)(\xi - u_2)(\xi - u_3)}} = \xi^{n-1/2} + \mathcal{O}(\xi^{-3/2}) \quad \text{for large } |\xi|$$

and

$$(2.3) \quad \int_{u_3}^{u_2} \frac{P_n(\xi, u_1, u_2, u_3)}{\sqrt{(\xi - u_1)(\xi - u_2)(\xi - u_3)}} d\xi = 0 .$$

Here the sign of the square root is given by $\sqrt{(\xi - u_1)(\xi - u_2)(\xi - u_3)} > 0$ for $\xi > u_1$ and the branch cuts are along $(-\infty, u_3)$ and (u_2, u_1) .

In particular,

$$(2.4) \quad P_0(\xi, u_1, u_2, u_3) = \xi + a_{0,1} , \quad P_1(\xi, u_1, u_2, u_3) = \xi^2 - \frac{1}{2}(u_1 + u_2 + u_3)\xi + a_{1,2} ,$$

where

$$a_{0,1} = (u_1 - u_3) \frac{E(s)}{K(s)} - u_1 ,$$

$$a_{1,2} = \frac{1}{3}(u_1 u_2 + u_1 u_3 + u_2 u_3) + \frac{1}{6}(u_1 + u_2 + u_3) a_{0,1} .$$

Here

$$s = \frac{u_2 - u_3}{u_1 - u_3}$$

and $K(s)$ and $E(s)$ are complete elliptic integrals of the first and second kind.

$K(s)$ and $E(s)$ have some well-known properties [10, 11]. They have the expansions

$$(2.5) \quad K(s) = \frac{\pi}{2} \left[1 + \frac{s}{4} + \frac{9}{64}s^2 + \dots + \left(\frac{1 \cdot 3 \cdots (2n-1)}{2 \cdot 4 \cdots 2n} \right)^2 s^n + \dots \right] ,$$

$$(2.6) \quad E(s) = \frac{\pi}{2} \left[1 - \frac{s}{4} - \frac{3}{64}s^2 - \dots - \frac{1}{2n-1} \left(\frac{1 \cdot 3 \cdots (2n-1)}{2 \cdot 4 \cdots 2n} \right)^2 s^n - \dots \right] ,$$

for $|s| < 1$. They also have the asymptotics

$$(2.7) \quad K(s) \approx \frac{1}{2} \log \frac{16}{1-s} ,$$

$$(2.8) \quad E(s) \approx 1 + \frac{1}{4}(1-s) \left[\log \frac{16}{1-s} - 1 \right] ,$$

as s is close to 1. Furthermore,

$$(2.9) \quad \frac{dK(s)}{ds} = \frac{E(s) - (1-s)K(s)}{2s(1-s)},$$

$$(2.10) \quad \frac{dE(s)}{ds} = \frac{E(s) - K(s)}{2s}.$$

It immediately follows from (2.5) and (2.6) that

$$(2.11) \quad \frac{1}{1 - \frac{s}{2}} < \frac{K(s)}{E(s)} < \frac{1 - \frac{s}{2}}{1 - s} \quad \text{for } 0 < s < 1.$$

The eigenspeeds of the Whitham equations (1.3) are defined in terms of P_0 and P_1 of (2.4),

$$\lambda_i(u_1, u_2, u_3) = 12 \frac{P_1(u_i, u_1, u_2, u_3)}{P_0(u_i, u_1, u_2, u_3)}, \quad i = 1, 2, 3,$$

which give

$$(2.12) \quad \begin{aligned} \lambda_1(u_1, u_2, u_3) &= 2(u_1 + u_2 + u_3) + 4(u_1 - u_2) \frac{K(s)}{E(s)}, \\ \lambda_2(u_1, u_2, u_3) &= 2(u_1 + u_2 + u_3) + 4(u_2 - u_1) \frac{sK(s)}{E(s) - (1-s)K(s)}, \\ \lambda_3(u_1, u_2, u_3) &= 2(u_1 + u_2 + u_3) + 4(u_2 - u_3) \frac{K(s)}{E(s) - K(s)}. \end{aligned}$$

In view of (2.5-2.8), we find that λ_1 , λ_2 and λ_3 have behavior:

(1) At $u_2 = u_3$:

$$(2.13) \quad \begin{aligned} \lambda_1(u_1, u_2, u_3) &= 6u_1, \\ \lambda_2(u_1, u_2, u_3) &= \lambda_3(u_1, u_2, u_3) = 12u_3 - 6u_1. \end{aligned}$$

(2) At $u_1 = u_2$:

$$(2.14) \quad \begin{aligned} \lambda_1(u_1, u_2, u_3) &= \lambda_2(u_1, u_2, u_3) = 4u_1 + 2u_3, \\ \lambda_3(u_1, u_2, u_3) &= 6u_3. \end{aligned}$$

The eigenspeeds of the Whitham equations (1.6) are

$$(2.15) \quad \mu_i^{(m)}(u_1, u_2, u_3) = 4^m (2m + 1) \frac{P_m(u_i, u_1, u_2, u_3)}{P_0(u_i, u_1, u_2, u_3)}, \quad i = 1, 2, 3.$$

The polynomial $4^m(2m+1)P_m(\xi, u_1, u_2, u_3)$ can be expressed as [2]

$$(2.16) \quad 4^m(2m+1)P_m(\xi, u_1, u_2, u_3) = 2(\xi - u_1)(\xi - u_2)(\xi - u_3)\Phi(\xi, u_1, u_2, u_3) + Q(\xi, u_1, u_2, u_3) .$$

The function $\Phi(\xi, \vec{u})$ satisfies the boundary value problem for the Euler-Poisson-Darboux equations

$$(2.17) \quad 2(u_i - u_j) \frac{\partial^2 \Phi}{\partial u_i \partial u_j} = \frac{\partial \Phi}{\partial u_i} - \frac{\partial \Phi}{\partial u_j} ,$$

$$(2.18) \quad 2(\xi - u_i) \frac{\partial^2 \Phi}{\partial \xi \partial u_i} = \frac{\partial \Phi}{\partial \xi} - 2 \frac{\partial \Phi}{\partial u_i} ,$$

$$(2.19) \quad \Phi(u, u, u, u) = \frac{2}{3} \frac{d^2}{du^2} [C_m u^m] ,$$

where

$$(2.20) \quad C_m = \frac{2^{2m+1}}{\int_0^1 \frac{t^m}{\sqrt{1-t}} dt} = \frac{2^m(2m+1)!!}{m!} .$$

The function $Q(\xi, \vec{u})$ is a quadratic polynomial in ξ ;

$$(2.21) \quad \begin{aligned} Q(\xi, u_1, u_2, u_3) = & 2(\xi - u_2)(\xi - u_3) \frac{\partial q(u_1, u_2, u_3)}{\partial u_1} + 2(\xi - u_1)(\xi - u_3) \frac{\partial q(u_1, u_2, u_3)}{\partial u_2} \\ & + 2(\xi - u_1)(\xi - u_2) \frac{\partial q(u_1, u_2, u_3)}{\partial u_3} + q(u_1, u_2, u_3) P_0(\xi, u_1, u_2, u_3) \end{aligned}$$

and $q(\vec{u})$ is the solution of the boundary value problem for another version of the Euler-Poisson-Darboux equations

$$(2.22) \quad 2(u_i - u_j) \frac{\partial^2 q}{\partial u_i \partial u_j} = \frac{\partial q}{\partial u_i} - \frac{\partial q}{\partial u_j} , \quad i, j = 1, 2, 3 ,$$

$$(2.23) \quad q(u, u, u) = C_m u^m .$$

The solution of equations (2.17-2.19) and that of (2.22) and (2.23) can be solved explicitly [2]. In particular, the solution of (2.22) and (2.23) is [10]

$$(2.24) \quad q(u_1, u_2, u_3) = \frac{C_m}{2\sqrt{2}\pi} \int_{-1}^1 \int_{-1}^1 \frac{(\frac{1+\mu}{2} \frac{1+\nu}{2} u_1 + \frac{1+\mu}{2} \frac{1-\nu}{2} u_2 + \frac{1-\mu}{2} u_3)^m}{\sqrt{(1-\mu)(1-\nu^2)}} d\mu d\nu .$$

The speeds $\mu_i^{(m)}$'s of (2.15) for $m > 1$ are connected to $\mu_i^{(m)}$'s for $m = 1$, which are also given by (2.12).

Lemma 2.1. For $i = 1, 2, 3$,

$$(2.25) \quad \mu_i^{(m)}(u_1, u_2, u_3) = \frac{1}{2}[\lambda_i(u_1, u_2, u_3) - 2(u_1 + u_2 + u_3)] \frac{\partial q(u_1, u_2, u_3)}{\partial u_i} + q(u_1, u_2, u_3) .$$

Proof. We use (2.15), (2.16) and (2.21) to write

$$(2.26) \quad \begin{aligned} \mu_1^{(m)}(u_1, u_2, u_3) &= \frac{Q(u_1, u_1, u_2, u_3)}{P_0(u_1, u_1, u_2, u_3)} \\ &= \frac{2(u_1 - u_2)(u_1 - u_3)}{P_0(u_1, u_1, u_2, u_3)} \frac{\partial q(u_1, u_2, u_3)}{\partial u_i} + q(u_1, u_2, u_3) . \end{aligned}$$

In particular, when $m = 1$, since the corresponding $q = 2(u_1 + u_2 + u_3)$, we obtain

$$\lambda_1(u_1, u_2, u_3) = \frac{4(u_1 - u_2)(u_1 - u_3)}{P(u_1, u_1, u_2, u_3)} + 2(u_1 + u_2 + u_3) .$$

This together with (2.26) proves formula (2.25) for $i = 1$. The cases for $i = 2, 3$ can be shown in the same way. □

Lemma 2.2. [12]

1.

$$(2.27) \quad \frac{\partial \mu_i^{(m)}}{\partial u_j} = \frac{\frac{\partial \lambda_i}{\partial u_j}}{\lambda_i - \lambda_j} [\mu_i^{(m)} - \mu_j^{(m)}] , \quad i, j = 1, 2, 3; i \neq j .$$

2.

$$(2.28) \quad \frac{\partial}{\partial u_i} \left(\frac{P_m(\xi, u_1, u_2, u_3)}{\sqrt{(\xi - u_1)(\xi - u_2)(\xi - u_3)}} \right) = \frac{\mu_i^{(m)}(u_1, u_2, u_3)}{4^m(2m+1)} \frac{\partial}{\partial u_i} \left(\frac{P_0(\xi, u_1, u_2, u_3)}{\sqrt{(\xi - u_1)(\xi - u_2)(\xi - u_3)}} \right)$$

for $i = 1, 2, 3$ and $\eta \neq u_1, u_2, u_3$.

The following calculations are useful in the subsequent sections.

Using formula (2.25) for μ_2 and μ_3 and formulae (2.12) for λ_2 and λ_3 , we obtain

$$(2.29) \quad \mu_2^{(m)}(u_1, u_2, u_3) - \mu_3^{(m)}(u_1, u_2, u_3) = \frac{2(u_2 - u_3)K}{(K - E)[E - (1 - s)K]} M(u_1, u_2, u_3) ,$$

where

$$(2.30) \quad M(u_1, u_2, u_3) = \left[\frac{\partial q}{\partial u_3} + (1-s) \frac{\partial q}{\partial u_2} \right] E - (1-s) \left(\frac{\partial q}{\partial u_2} + \frac{\partial q}{\partial u_3} \right) K .$$

We then use (2.9), (2.10) and (2.22) to calculate

$$(2.31) \quad \frac{\partial M(u_1, u_2, u_3)}{\partial u_2} = \frac{1}{2} \frac{p_1(u_1, u_2, u_3)}{u_1 - u_3} [E - K] ,$$

$$(2.32) \quad \frac{\partial M(u_1, u_2, u_3)}{\partial u_3} = \frac{1}{2} \frac{p_2(u_1, u_2, u_3)}{u_1 - u_3} [E - (1-s)K] + \frac{3}{2} \frac{M(u_1, u_2, u_3)}{u_1 - u_3} ,$$

where

$$(2.33) \quad p_1(u_1, u_2, u_3) = 2(u_1 - u_2) \frac{\partial}{\partial u_2} \operatorname{div}(q) - \operatorname{div}(q) , \quad p_2(u_1, u_2, u_3) = 2(u_1 - u_3) \frac{\partial}{\partial u_3} \operatorname{div}(q) - \operatorname{div}(q) .$$

We next consider

$$(2.34) \quad F(u_1, u_2, u_3) := \frac{\mu_2(u_1, u_2, u_3) - \mu_3(u_1, u_2, u_3)}{u_2 - u_3} .$$

Using formula (2.25) for μ_2 and μ_3 and formulae (2.12) for λ_2 and λ_3 , we obtain

$$\begin{aligned} F &= -2 \frac{(1-s)K}{E - (1-s)K} \frac{\partial q}{\partial u_2} + 2 \frac{K}{K - E} \frac{\partial q}{\partial u_3} \\ &= -4 \frac{s(1-s)K}{E - (1-s)K} (u_1 - u_3) \frac{\partial^2 q}{\partial u_2 \partial u_3} + 2 \left[\frac{K}{K - E} - \frac{(1-s)K}{E - (1-s)K} \right] \frac{\partial q}{\partial u_3} , \end{aligned}$$

where we have used equations (2.22) in the last equality. Finally, we use the expansions (2.5-2.6) for K and E to obtain

$$(2.35) \quad F(u_1, u_2, u_3) = -4 \left[\left(2 - \frac{7}{4}s + \cdots \right) (u_1 - u_3) \frac{\partial^2 q}{\partial u_2 \partial u_3} + \left(-\frac{3}{4} + O(s^2) \right) \frac{\partial q}{\partial u_3} \right] .$$

3. SELF-SIMILAR SOLUTIONS

In this section, we construct the self-similar solution of the Whitham equations (1.6) when $m \geq 2$ for the initial function (1.7). The $m = 2$ result has already been obtained in [9]. Even in the $m = 2$ case, the key calculations presented here are different from those in [9].

Theorem 3.1. (see Figure 1.) For the step-like initial data $u_0(x)$ of (1.7), the solution of the Whitham equations (1.6) is given by

$$(3.1) \quad u_1 = 1, \quad x = \mu_2^{(m)}(1, u_2, u_3) t, \quad x = \mu_3^{(m)}(1, u_2, u_3) t$$

for $\alpha t < x \leq \beta t$ and by

$$(3.2) \quad u_1 = 1, \quad x = \mu_2^{(m)}(1, u_2, 0) t, \quad u_3 = 0$$

for $\beta t \leq x < \gamma t$, where $\alpha = \mu_2^{(m)}(1, u^*, u^*)$, $\beta = \mu_2^{(m)}(1, u^{**}, 0)$ and $\gamma = \mu_2^{(m)}(1, 1, 0) = q(1, 1, 0) = 4^m$.

Here, u^* is uniquely determined by the equation

$$(3.3) \quad p_1(1, u^*, u^*) = 0,$$

and u^{**} is uniquely given by the equation

$$(3.4) \quad \mu_2^{(m)}(1, u^{**}, 0) - \mu_3^{(m)}(1, u^{**}, 0) = 0.$$

Outside the region $\alpha t < x < 4^m t$, the solution of the Burgers type equation (1.5) is given by

$$(3.5) \quad u \equiv 1 \quad x \leq \alpha t$$

and

$$(3.6) \quad u \equiv 0 \quad x \geq 4^m t.$$

The boundaries $x = \alpha t$ and $x = 4^m t$ are called the trailing and leading edges, respectively. They separate the solutions of the Whitham equations and Burgers type equations. The Whitham solution matches the Burgers type solution in the following fashion (see Figure 1.):

$$(3.7) \quad u_1 = \text{the Burgers type solution defined outside the region},$$

$$(3.8) \quad u_2 = u_3,$$

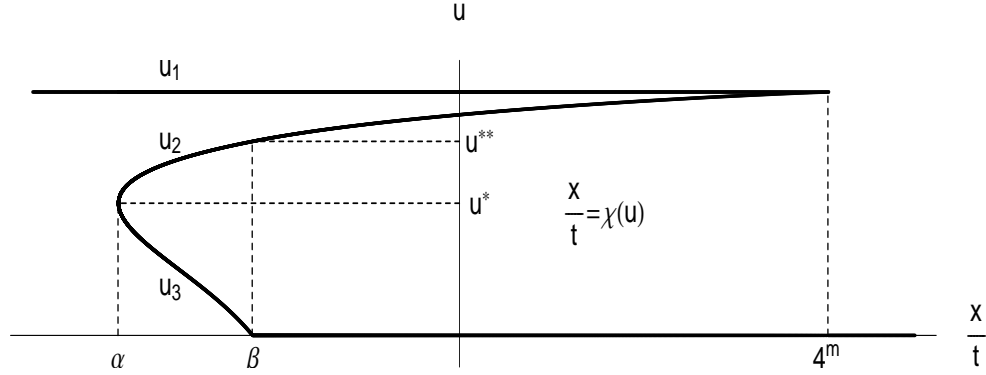


FIGURE 1. Self-Similar solution of the Whitham equations for $m \geq 2$. The curve defines the function $\chi(u)$.

at the trailing edge;

$$(3.9) \quad u_1 = u_2 ,$$

$$(3.10) \quad u_3 = \text{the Burgers type solution defined outside the region ,}$$

at the leading edge.

The proof of Theorem 3.1 is based on a series of lemmas.

We first show that the solution defined by either formulae (3.1) or (3.2) indeed satisfies the Whitham equations (1.6) [1, 9, 13].

Lemma 3.2. (1) *The functions u_1 , u_2 and u_3 determined by equations (3.1) give a solution of the Whitham equations (1.6) as long as u_2 and u_3 can be solved from (3.1) as functions of x and t .*

(2) *The functions u_1 , u_2 and u_3 determined by equations (3.2) give a solution of the Whitham equations (1.6) as long as u_2 can be solved from (3.2) as a function of x and t .*

Proof. (1) u_1 obviously satisfies the first equation of (1.6). To verify the second and third equations, we observe that

$$(3.11) \quad \frac{\partial \mu_2^{(m)}}{\partial u_3} = \frac{\partial \mu_3^{(m)}}{\partial u_2} = 0$$

on the solution of (3.1). To see this, we use (2.27) to calculate

$$\frac{\partial \mu_2^{(m)}}{\partial u_3} = \frac{\frac{\partial \lambda_2}{\partial u_3}}{\lambda_2 - \lambda_3} (\mu_2^{(m)} - \mu_3^{(m)}) = 0 .$$

The second part of (3.11) can be shown in the same way.

We then calculate the partial derivatives of the second equation of (3.1) with respect to x and t .

$$1 = \frac{\partial \mu_2^{(m)}}{\partial u_2} t u_{2x} , \quad 0 = \frac{\partial \mu_2^{(m)}}{\partial u_2} t u_{2t} + \mu_2^{(m)} ,$$

which give the second equation of (1.6).

The third equation of (1.6) can be verified in the same way.

(2) The second part of Lemma 3.2 can easily be proved.

□

We now determine the trailing edge. Eliminating x and t from the last two equations of (3.1) yields

$$(3.12) \quad \mu_2^{(m)}(1, u_2, u_3) - \mu_3^{(m)}(1, u_2, u_3) = 0 .$$

Since it degenerates at $u_2 = u_3$, we replace (3.12) by

$$(3.13) \quad F(1, u_2, u_3) := \frac{\mu_2^{(m)}(1, u_2, u_3) - \mu_3^{(m)}(1, u_2, u_3)}{u_2 - u_3} = 0 .$$

Here, the function F is also defined in (2.34).

Therefore, at the trailing edge where $u_2 = u_3$, i.e., $s = 0$, equation (3.13), in view of the expansion (2.35), becomes

$$8(1 - u_2) \frac{\partial^2 q(1, u_2, u_2)}{\partial u_2 \partial u_3} - 3 \frac{\partial q(1, u_2, u_2)}{\partial u_3} = 0 .$$

Since $\frac{\partial q}{\partial u_2} = \frac{\partial q}{\partial u_3}$ and $\frac{\partial^2 q}{\partial u_2^2} = 3 \frac{\partial^2 q}{\partial u_2 \partial u_3}$ on $u_2 = u_3$ because of (2.22), this equation is exactly equation (3.3).

Lemma 3.3. *Equation $p_1(1, \xi, \xi) = 0$ has a simple zero, denoted by u^* , in the region $0 < \xi < 1$, counting multiplicities. Furthermore, $p_1(1, \xi, \xi)$ is positive when $\xi < u^*$ and negative when $\xi > u^*$.*

Proof. We first simplify the polynomial p_1 of (2.33). In view of formula (2.24) for q , we use the fact that q is symmetric in u_1 , u_2 and u_3 to obtain

$$(3.14) \quad \left(\frac{\partial q(u_1, u_2, u_3)}{\partial u_1} + \frac{\partial q(u_1, u_2, u_3)}{\partial u_2} + \frac{\partial q(u_1, u_2, u_3)}{\partial u_3} \right) \Big|_{u_2=u_3=\xi} = C_m U_0(\xi, u_1) ,$$

where

$$(3.15) \quad U_0(\xi, u_1) = \frac{1}{2\sqrt{2}} \int_{-1}^1 \frac{m(\frac{1+\mu}{2}\xi + \frac{1-\mu}{2}u_1)^{m-1}}{\sqrt{1-\mu}} d\mu .$$

We can then write p_1 as

$$(3.16) \quad p_1(1, \xi, \xi) = C_m [(1 - \xi) \frac{\partial U_0(\xi, 1)}{\partial \xi} - U_0(\xi, 1)] .$$

Denoting the function in the parenthesis of (3.16) by $p(\xi)$, we claim that

$$(3.17) \quad \frac{d^k p(0)}{d\xi^k} > 0 , \quad \frac{d^k p(1)}{d\xi^k} < 0$$

for $k = 0, 1, 2, \dots, m-2$.

Obviously,

$$\frac{d^k p(\xi)}{d\xi^k} = (1 - \xi) \frac{d^{k+1} U_0(\xi, 1)}{d\xi^{k+1}} - (k+1) \frac{d^k U_0(\xi, 1)}{d\xi^k} .$$

Since, $\frac{d^k U_0(\xi, 1)}{d\xi^k}$ is a positive function, this proves the second inequality of (3.17).

To prove the first inequality of (3.17), we use formula (3.15) to calculate

$$\frac{d^k U_0(\xi, 1)}{d\xi^k} \Big|_{\xi=0} = \frac{m(m-1) \cdots (m-k)}{2^{m+\frac{1}{2}}} \int_{-1}^1 (1-\mu)^{m-k-\frac{3}{2}} (1+\mu)^k d\mu .$$

The integral on the right can be evaluated using an iteration formula. Denote this integral by $A_{m,k}$.

An integration by parts gives $A_{m,k} = \frac{2k}{2m-2k-1} A_{m,k-1}$. Since $A_{m,0} = 2^{m+\frac{1}{2}}/(2m-1)$, we thus obtain

$$A_{m,k} = \frac{2^{m+k+\frac{1}{2}} k!}{(2m-1)(2m-3) \cdots (2m-2k-1)} ,$$

which gives

$$(3.18) \quad \frac{d^k U_0(\xi, 1)}{d\xi^k} \Big|_{\xi=0} = \frac{2^k k! m(m-1) \cdots (m-k)}{(2m-1)(2m-3) \cdots (2m-2k-1)} .$$

Therefore

$$\frac{d^k p(0)}{d\xi^k} = \frac{d^{k+1} U_0(\xi, 1)}{d\xi^{k+1}} \Big|_{\xi=0} - (k+1) \frac{d^k U_0(\xi, 1)}{d\xi^k} \Big|_{\xi=0} = \frac{2^k (k+1)! m(m-1) \cdots (m-k)}{(2m-1)(2m-3) \cdots (2m-2k-3)} > 0 .$$

We now use (3.17) to prove the existence and uniqueness of the zero of function $p^{(k)}(\xi)$, $k = 0, 1, \dots, (m-2)$. First, it follows from (3.17) that $p^{(k)}(\xi)$ has an odd number of zeros in $0 < \xi < 1$, counting multiplicities. Second, if $p^{(k)}(\xi)$ has more than one zero, it must have at least three zeros. Consequently, $p^{(k+1)}(\xi)$ will have more than one zero; so it must also have at least three zeros. Repeating this argument, we see that $p^{(m-2)}$ must have at least three zeros. This is an impossibility since $p(\xi)$ is a polynomial of degree $m-1$ because $U_0(\xi, 1)$ is so. Therefore, $p^{(k)}(\xi)$ has one and only one zero for $\xi \in (0, 1)$ when $k = 0, 1, \dots, (m-2)$. In particular, the $k = 0$ case proves Lemma 3.3. \square

Lemma 3.4. *Equation (3.13) has a unique solution satisfying $u_2 = u_3$. The solution is $u_2 = u_3 = u^*$. The rest of equations (3.1) at the trailing edge are $u_1 = 1$ and $x/t = \mu_2^{(m)}(1, u^*, u^*)$.*

Having located the trailing edge, we now solve equations (3.1) in the neighborhood of the trailing edge. We first consider equation (3.13). We use (2.35) to differentiate F at the trailing edge $u_1 = 1$, $u_2 = u_3 = u^*$

$$\begin{aligned} \frac{\partial F(1, u^*, u^*)}{\partial u_2} &= \frac{\partial F(1, u^*, u^*)}{\partial u_3} = 10 \frac{\partial^2 q(1, u^*, u^*)}{\partial u_2 \partial u_3} - 8(1 - u^*) \frac{\partial^3 q(1, u^*, u^*)}{\partial u_2^2 \partial u_3} > 0 \\ (3.19) \qquad \qquad \qquad &= -\frac{C_m}{2} \frac{\partial}{\partial \xi} \left[(1 - \xi) \frac{\partial U_0(\xi, 1)}{\partial \xi} - U_0(\xi, 1) \right]_{\xi=u^*} > 0 , \end{aligned}$$

where in the second equality we have used (2.22), (3.14) and identities $\frac{\partial^2 q}{\partial u_2^2} = \frac{\partial^2 q}{\partial u_3^2} = 3 \frac{\partial^2 q}{\partial u_2 \partial u_3}$ on $u_2 = u_3$. The inequality is a consequence of Lemma 3.3.

Inequality (3.19) shows that equation (3.13) or equivalently (3.12) can be inverted to give u_2 as a decreasing function of u_3

$$(3.20) \qquad \qquad \qquad u_2 = B(u_3)$$

in a neighborhood of $u_2 = u_3 = u^*$.

We will extend the solution (3.20) of equation (3.12) by decreasing u_3 in the region $0 < u_3 < u^* < u_2 < 1$ as far as possible. We need to evaluate the derivatives $\frac{\partial(\mu_2^{(m)} - \mu_3^{(m)})}{\partial u_2}$ and $\frac{\partial(\mu_2^{(m)} - \mu_3^{(m)})}{\partial u_3}$ on the solution of (3.1). It follows from (2.29), (2.31) and (2.32) that

$$(3.21) \quad \frac{\partial[\mu_2^{(m)} - \mu_3^{(m)}]}{\partial u_2} = -\frac{sKp_1(1, u_2, u_3)}{E - (1-s)K},$$

$$(3.22) \quad \frac{\partial[\mu_2^{(m)} - \mu_3^{(m)}]}{\partial u_3} = \frac{sKp_2(1, u_2, u_3)}{K - E},$$

on the solution of (3.1).

We first study the two polynomials p_1 and p_2 of (2.33).

Lemma 3.5. *For each $0 \leq u_3 < 1$, the polynomial $p_1(1, u_2, u_3)$, as a function of u_2 , has only one zero in the region $0 < u_2 < 1$, counting multiplicities. Furthermore, $p_1(1, u_2, u_3)$ is positive when u_2 is on the left of this zero and negative when u_2 is on the right.*

For each $0 \leq u_2 < 1$, the polynomial $p_2(1, u_2, u_3)$, as a function of u_3 , has only one zero in the region $0 < u_3 < 1$, counting multiplicities. Furthermore, $p_2(1, u_2, u_3)$ is positive when u_3 is on the left of this zero and negative when u_3 is on the right.

Proof. We will prove the first part of the lemma; the second part follows from $p_2(1, u_2, u_3) = p_1(1, u_3, u_2)$.

The proof of the first part is similar to the proof of Lemma 3.3. We will go through it briefly.

We first have

$$\frac{\partial^k p_1(1, 0, u_3)}{\partial u_2^k} > 0, \quad \frac{\partial^k p_1(1, 1, u_3)}{\partial u_2^k} < 0$$

for $k = 0, 1, 2, \dots, m-2$. The second inequality immediately follows from formula (2.33) for p_1 . The first inequality is derived from a formula similar to (3.18).

The rest of the proof is the same as the proof of Lemma 3.3. □

We now continue to extend the solution (3.20) of equation (3.12) in the region $0 < u_3 < u^* < u_2 < 1$ as far as possible. When u_2 and u_3 are close to u^* , because of (3.19), we have $\frac{\partial(\mu_2^{(m)} - \mu_3^{(m)})}{\partial u_2} > 0$ and $\frac{\partial(\mu_2^{(m)} - \mu_3^{(m)})}{\partial u_3} > 0$ on the solution of (3.12). These along with (3.21) and (3.22) show that

$$(3.23) \quad p_1(1, u_2, u_3) < 0, \quad p_2(1, u_2, u_3) > 0$$

when u_2 and u_3 are close to u^* .

Lemma 3.6. *Inequalities (3.23) hold on the solution of equation (3.12) as long as $0 < u_3 < u^* < u_2 < 1$.*

Proof. We first prove the first inequality of (3.23). In view of (2.29), the solution of equation (3.12) is also governed by an equivalent equation $M(1, u_2, u_3) = 0$. For each $0 < u_3 < u^*$, we study the zero of the u_2 -variable function $M(1, u_2, u_3)$ in the interval $u_3 < u_2 < 1$. We obtain from formula (2.30) for M that $M(1, u_3, u_3) = 0$ and $M(1, 1, u_3) > 0$. $M(1, u_2, u_3)$ is a decreasing function of u_2 when u_2 is on the immediate right of u_3 . To see this, we note that $p_1(1, u_3, u_3) > 0$ for $u_3 < u^*$ according to Lemma 3.3. This and (2.31) prove that M is decreasing for u_2 on the immediate right of u_3 because $E - K < 0$ for $s > 0$; so $M(1, u_2, u_3) < 0$ for such u_2 . Therefore, $M(1, u_2, u_3)$ has a u_2 -zero in the interval $u_3 < u_2 < 1$ when $0 < u_3 < u^*$. Because of the uniqueness of the u_2 -zero of $p(1, u_2, u_3)$ according to Lemma 3.5, we conclude that $M(1, u_2, u_3)$ has only one zero and that this zero is on the right of the zero of p_1 . Hence, the zero of M is exactly given by $u_2 = B(u_3)$ and p_1 is negative on the solution. This proves the first inequality of (3.23).

We now prove the second inequality of (3.23) by contradiction. Suppose it first fails at \bar{u}_2 and \bar{u}_3 , where $0 < \bar{u}_3 < u^* < \bar{u}_2 < 1$; i.e.,

$$(3.24) \quad p_2(1, B(u_3), u_3) > 0 \text{ when } \bar{u}_3 < u_3 < u^*, \quad p_2(1, B(u_3), u_3) = 0 \text{ when } u_3 = \bar{u}_3.$$

It then follows from (3.21) and (3.22) that solution (3.20) of equation (3.12) has a zero derivative at $u_3 = \bar{u}_3$; i.e.,

$$B'(\bar{u}_3) = 0 .$$

Hence,

$$\frac{d}{du_3} p_2(1, B(u_3), u_3)|_{u_3=\bar{u}_3} = \frac{\partial}{\partial u_2} p_2(1, B(\bar{u}_3), \bar{u}_3) B'(\bar{u}_3) + \frac{\partial}{\partial u_3} p_2(1, B(\bar{u}_3), \bar{u}_3) < 0 ,$$

where the first term vanishes because of $B'(\bar{u}_3) = 0$ and the second term is negative according to Lemma 3.5. In view of $p_2(1, B(\bar{u}_3), \bar{u}_3) = 0$, this implies that $p_2(1, B(u_3), u_3) < 0$ when u_3 is on the immediately right of \bar{u}_3 . That contradicts (3.24). This proves the second inequality of (3.23). \square

It follows from (3.21), (3.22) and Lemma 3.6 that

$$(3.25) \quad \frac{\partial[\mu_2^{(m)} - \mu_3^{(m)}]}{\partial u_2} > 0 , \quad \frac{\partial[\mu_2^{(m)} - \mu_3^{(m)}]}{\partial u_3} > 0$$

on the solution of (3.12). Solution (3.20) of equation (3.12) can then be extended as a decreasing function of u_3 as long as $0 < u_3 < u^* < u_2 < 1$.

There are two possibilities: (1) u_2 touches 1 before or simultaneously as u_3 reaches 0 and (2) u_3 touches 0 before u_2 reaches 1.

It follows from (2.14) and (2.25) that

$$\mu_2^{(m)}(1, 1, u_3) > \mu_3^{(m)}(1, 1, u_3) \quad \text{for } 0 \leq u_3 < 1 .$$

This shows that (1) is impossible. Hence, u_3 will touch 0 before u_2 reaches 1. When this happens, equation (3.12) becomes equation (3.4).

Lemma 3.7. *Equation (3.4) has a simple zero in the region $0 < u_2 < 1$, counting multiplicities. Denoting the zero by u^{**} , then $\mu_2^{(m)}(1, u_2, 0) - \mu_3^{(m)}(1, u_2, 0)$ is positive for $u_2 > u^{**}$ and negative for $u_2 < u^{**}$. Furthermore, $0 < u^{***} < u^{**} < 1$ where u^{***} is the unique zero of $p_1(1, u_2, 0)$.*

Proof. We use (2.29) and (2.31) to prove the lemma. In equation (2.29), $K - E$ and $E - (1 - s)K$ are all positive for $0 < s < 1$ in view of (2.11).

Denoting the unique zero of $p_1(1, u_2, 0)$ by u^{***} , it then follows from Lemma 3.5 that $p_1(1, u_2, 0) > 0$ when $0 < u_2 < u^{***}$ and $p_1(1, u_2, 0) < 0$ when $u^{***} < u_2 < 1$. Since $M(1, u_2, 0)$ of (2.30) vanishes at $u_2 = 0$ and is positive at $u_2 = 1$ in view of (2.5-2.8), we conclude from the derivative (2.31) that $M(1, u_2, 0)$ has a simple zero in $0 < u_2 < 1$. This zero is exactly u^{**} and the rest of the theorem can be proved easily. \square

Having solved equation (3.12) for u_2 as a decreasing function of u_3 for $0 < u_3 < u^*$, we turn to equations (3.1). Because of (3.11) and (3.25), the third equation of (3.1) gives u_3 as a decreasing function of x/t for $\alpha \leq x/t \leq \beta$, where $\alpha = \mu_2^{(m)}(1, u^*, u^*)$ and $\beta = \mu_2^{(m)}(1, u^{**}, 0)$. Consequently, u_2 is an increasing function of x/t in the same interval.

Lemma 3.8. *The last two equations of (3.1) can be inverted to give u_2 and u_3 as increasing and decreasing functions, respectively, of the self-similarity variable x/t in the interval $\alpha \leq x/t \leq \beta$.*

We now turn to equations (3.2). We first use (2.9), (2.10) and (2.12) to calculate the derivative of $\mu_2^{(m)}$ of (2.25)

$$\begin{aligned}
 \frac{\partial \mu_2^{(m)}}{\partial u_2} &= \frac{1}{2} [\lambda_2 - 2(1 + u_2 + u_3)] \frac{\partial^2 q}{\partial u_2^2} + \frac{1}{2} \frac{\partial \lambda_2}{\partial u_2} \frac{\partial q}{\partial u_2} \\
 &= \frac{2(u_2 - 1)sK}{E - (1 - s)K} \frac{\partial^2 q}{\partial u_2^2} + \left[\frac{2sK}{E - (1 - s)K} + 1 - \frac{E^2 - (1 - s)K^2}{(E - (1 - s)K)^2} \right] \frac{\partial q}{\partial u_2} \\
 (3.26) \quad &> -\frac{2sK}{E - (1 - s)K} \left[(1 - u_2) \frac{\partial^2 q}{\partial u_2^2} - \frac{\partial q}{\partial u_2} \right],
 \end{aligned}$$

where in the inequality we have used $(E - (1 - s)K)^2 > E^2 - (1 - s)K^2$, which is a consequence of (2.11).

The polynomial in the parenthesis of (3.26) is connected to $p_1(1, u_2, 0)$; indeed,

$$(3.27) \quad p_1(1, u_2, 0) = \frac{2m + 1}{m} \left[(1 - u_2) \frac{\partial^2 q(1, u_2, 0)}{\partial u_2^2} - \frac{\partial q(1, u_2, 0)}{\partial u_2} \right].$$

This follows from the identity

$$(3.28) \quad 2m \frac{\partial q(1, u_2, 0)}{\partial u_3} = (1 - u_2) \frac{\partial q(1, u_2, 0)}{\partial u_2} + mq(1, u_2, 0) .$$

To see this, taking the derivative of (3.28) and using formula (2.33) for p_1 and equations (2.22) for q yield (3.27).

To prove (3.28), we use the integral formula (2.24) for q to calculate both sides of the identity.

The left equals

$$\frac{C_m m^2}{4^{m-1} 2^{3/2} \pi} \left[\int_{-1}^1 (1 + \mu)^{m-1} (1 - \mu)^{1/2} d\mu \right] \left[\int_{-1}^1 ((1 + \nu) + (1 - \nu)u_2)^{m-1} (1 - \nu^2)^{-1/2} d\nu \right] .$$

The right is

$$\frac{C_m m}{4^{m-1} 2^{5/2} \pi} \left[\int_{-1}^1 (1 + \mu)^m (1 - \mu)^{-1/2} d\mu \right] \left[\int_{-1}^1 ((1 + \nu) + (1 - \nu)u_2)^{m-1} (1 - \nu^2)^{-1/2} d\nu \right] .$$

Both sides are equal in view of an easy identity

$$m \int_{-1}^1 (1 + \mu)^{m-1} (1 - \mu)^{1/2} d\mu = \frac{1}{2} \int_{-1}^1 (1 + \mu)^m (1 - \mu)^{-1/2} d\mu .$$

We have therefore proved identity (3.28).

By Lemma 3.5, $p_1(1, u_2, 0)$ is negative for $u_2 > u^{***}$, where u^{***} is the unique zero of p_1 . Since $u^{***} < u^{**}$ according to Lemma 3.7, we conclude from (3.26) and (3.27) that $\frac{\partial \mu_2^{(m)}}{\partial u_2} > 0$ on the solution of (3.2) when $u_2 > u^{**}$. Hence, the second equation of (3.2) can be solved for u_2 as an increasing function of x/t as long as $u^{**} < u_2 < 1$. When u_2 reaches 1, we have

$$x/t = \gamma = \mu_2^{(m)}(1, 1, 0) .$$

We have therefore proved the following result.

Lemma 3.9. *The second equation of (3.2) can be inverted to give u_2 as an increasing function of x/t in the interval $\beta \leq x/t \leq \gamma$.*

We are ready to conclude the proof of Theorem 3.1.

The Burgers type solutions (3.5) and (3.6) are trivial.

According to Lemma 3.8, the last two equations of (3.1) determine u_2 and u_3 as functions of x/t in the region $\alpha \leq x/t \leq \beta$. By the first part of Lemma 3.2, the resulting u_1 , u_2 and u_3 satisfy the Whitham equations (1.6). Furthermore, the boundary conditions (3.7) and (3.8) are satisfied at the trailing edge $x = \alpha t$.

Similarly, by Lemma 3.9, the second equation of (3.2) determines u_2 as a function of x/t in the region $\beta \leq x/t \leq \gamma = 4^m$. It then follows from the second part of Lemma 3.2 that u_1 , u_2 and u_3 of (3.2) satisfy the Whitham equations (1.6). They also satisfy the boundary conditions (3.9) and (3.10) at the leading edge $x = \gamma t$.

We have therefore completed the proof of Theorem 3.1.

4. THE MINIMIZATION PROBLEM

The zero dispersion limit of the solution of the higher order KdV equation (1.4) with step-like initial function (1.7) is also determined by a minimization problem with constraints [5, 6, 14]

$$(4.1) \quad \text{Minimize}_{\{\psi \geq 0, \psi \in L^1\}} \left\{ -\frac{1}{2\pi} \int_0^1 \int_0^1 \log \left| \frac{\eta - \mu}{\eta + \mu} \right| \psi(\eta) \psi(\mu) d\eta d\mu + \int_0^1 [\eta x - 4^m \eta^{2m+1} t] \psi(\eta) d\eta \right\} .$$

In this section, we will use the self-similar solution of Section 3 to construct the minimizer for $m \geq 2$. The $m = 2$ result has already been obtained in [9]. Even in the $m = 2$ case, the key calculations presented here are different from those in [9].

We first define a linear operator

$$L\psi(\eta) = \frac{1}{2\pi} \int_0^1 \log \left(\frac{\eta - \mu}{\eta + \mu} \right)^2 \psi(\mu) d\mu .$$

The variational conditions are

$$(4.2) \quad L\psi = x\eta - 4^m t \eta^{2m+1} \quad \text{where } \psi > 0 ,$$

$$(4.3) \quad L\psi \leq x\eta - 4^m t \eta^{2m+1} \quad \text{where } \psi = 0 .$$

The constraint for the minimization problem is

$$(4.4) \quad \psi \geq 0 .$$

The minimizer of (4.1) is given explicitly:

Theorem 4.1. *The minimizer of the variational problem (4.1) is as follows:*

(1) For $x \leq \alpha t$,

$$\psi(\eta) = \frac{-x\eta + 4^m(2m+1)t\eta \frac{P_m(\eta^2, 1, u^*, u^*)}{\eta^2 - u^*}}{\sqrt{1 - \eta^2}} .$$

(2) For $\alpha t < x < \beta t$,

$$\psi(\eta) = \begin{cases} -\frac{-x\eta P_0(\eta^2, 1, u_2, u_3) + 4^m(2m+1)t\eta P_m(\eta^2, 1, u_2, u_3)}{\sqrt{(1-\eta^2)(u_2-\eta^2)(u_3-\eta^2)}} & 0 < \eta < \sqrt{u_3} \\ 0 & \sqrt{u_3} < \eta < \sqrt{u_2} \\ \frac{-x\eta P_0(\eta^2, 1, u_2, u_3) + 4^m(2m+1)t\eta P_m(\eta^2, 1, u_2, u_3)}{\sqrt{(1-\eta^2)(\eta^2-u_2)(\eta^2-u_3)}} & \sqrt{u_2} < \eta < 1 , \end{cases}$$

where P_0 and P_m are defined in (2.1) and u_2 and u_3 are determined by equations (3.1).

(3) For $\beta t < x < 4^m t$,

$$\psi(\eta) = \begin{cases} 0 & 0 < \eta < \sqrt{u_2} \\ \frac{-xP_0(\eta^2, 1, u_2, 0) + 4^m(2m+1)tP_m(\eta^2, 1, u_2, 0)}{\sqrt{(1-\eta^2)(\eta^2-u_2)}} & \sqrt{u_2} < \eta < 1 , \end{cases}$$

where u_2 is determined by (3.2).

(4) For $x \geq 4^m t$,

$$\psi(\eta) \equiv 0 .$$

Proof. We extend the function ψ defined on $[0, 1]$ to the entire real line by setting $\psi(\eta) = 0$ for $\eta > 1$ and taking ψ to be odd. In this way, the operator L is connected to the Hilbert transform H on the real line [5]:

$$L\psi(\eta) = \int_0^\eta H\psi(\mu) d\mu \quad \text{where} \quad H\psi(\eta) = \frac{1}{\pi} P.V. \int_{-\infty}^{+\infty} \frac{\psi(\mu)}{\eta - \mu} d\mu .$$

We verify case (4) first. Clearly $\psi(\eta) = 0$ satisfies the constraint (4.4). We now check the variational conditions (4.2-4.3). Since $\psi = 0$,

$$L\psi = 0 \leq x\eta - 4^m t \eta^{2m+1} ,$$

where the inequality follows from $x \geq 4^m t$ and $0 \leq \eta \leq 1$. Hence, variational conditions (4.2-4.3) are satisfied.

Next we consider case (1). We write $\psi(\eta)$ as the real part of $g_1(\eta)$ for real η , where

$$g_1 = \sqrt{-1}(x - 4^m(2m+1)t\eta^{2m}) + \frac{\sqrt{-1}[-x\eta + 4^m(2m+1)t\eta^{\frac{P_m(\eta^2, 1, u^*, u^*)}{\eta^2 - u^*}}]}{\sqrt{\eta^2 - 1}} .$$

The function g_1 is analytic in the upper half complex plane $Im(\eta) > 0$ and $g_1(\eta) \approx O(1/\eta^2)$ for large $|\eta|$ in view of the expansion (2.1) for P_m . Hence, $H\psi(\eta) = Im[g_1(\eta)] = x - 4^m(2m+1)t\eta^{2m}$ on $0 \leq \eta \leq 1$, where H is the Hilbert transform [5]. We then have for $0 \leq \eta \leq 1$

$$L\psi(\eta) = \int_0^\eta H\eta(\mu)d\mu = x\eta - 4^m t \eta^{2m+1} ,$$

which shows that the variational conditions (4.2) and (4.3) are satisfied.

To prove (4.4), we first claim that

$$(4.5) \quad -\alpha + 4^m(2m+1)\frac{P_m(\eta^2, 1, u^*, u^*)}{\eta^2 - u^*} \geq 0 ,$$

for $0 \leq \eta \leq 1$. To see this, we use (2.13) and (2.25) to calculate $\alpha = \mu_2^{(m)}(1, u^*, u^*)$ and (2.16) and (2.21) to evaluate $P_m(1, u^*, u^*)$. The left hand side of (4.5) equals

$$\begin{aligned} & 2(\eta^2 - 1)(\eta^2 - u^*)\Phi(\eta^2, 1, u^*, u^*) + 2(\eta^2 - u^*)\left[\frac{\partial q(1, u^*, u^*)}{\partial u_1} + \frac{\partial q(1, u^*, u^*)}{\partial u_2} + \frac{\partial q(1, u^*, u^*)}{\partial u_3}\right] \\ &= 2(\eta^2 - 1)[C_m U(\eta^2, 1) - C_m U(u^*, 1)] + 2(\eta^2 - u^*)C_m U(u^*, 1) \\ (4.6) \quad &= 2[(1 - u^*)C_m U(u^*, 1) - (1 - \eta^2)C_m U(\eta^2, 1)] , \end{aligned}$$

where in the first equality we have used (3.14) and the identity (cf. (3.28) of [2])

$$\Phi(\eta^2, 1, u^*, u^*) = \frac{C_m U(\eta^2, 1) - C_m U(u^*, 1)}{\eta^2 - u^*} .$$

In view of (3.16), we have

$$\frac{d}{d(\eta^2)}[(1 - \eta^2)C_m U(\eta^2, 1)] = C_m[(1 - \eta^2)\frac{\partial U(\eta^2, 1)}{\partial \eta^2} - U(\eta^2, 1)] = p_1(1, \eta^2, \eta^2) .$$

This derivative is positive when $\eta^2 < u^*$ and negative when $\eta^2 > u^*$ because p_1 is so according to Lemma 3.3. Therefore, $(1 - \eta^2)C_m U(\eta^2, 1)$ has a maximum at $\eta^2 = u^*$; i.e., $(1 - \eta^2)C_m U(\eta^2, 1) \leq (1 - u^*)C_m U(u^*, 1)$ for $0 \leq \eta^2 \leq 1$. This together with (4.6) proves (4.5).

It follows from $x \leq \alpha t$ and inequality (4.5) that $\psi \geq 0$. Hence, the constraint (4.4) is verified.

We now turn to case (2). By Lemma 3.5, the last two equations of (3.1) determine u_2 and u_3 as functions of the self-similarity variable x/t in the interval $\alpha \leq x/t \leq \beta$.

We write $\psi = \text{Re}(g_2)$ for real η , where

$$g_2 = \sqrt{-1}(x - 4^m(2m+1)t\eta^{2m}) + \frac{\sqrt{-1}[-x\eta P_0(\eta^2, 1, u_2, u_3) + 4^m(2m+1)t\eta P_m(\eta^2, 1, u_2, u_3)]}{\sqrt{\eta^2 - 1}(\eta^2 - u_2)(\eta^2 - u_3)} .$$

The function g_2 is analytic in $\text{Im}(\eta) > 0$ and $g_2(\eta) \approx O(1/\eta^2)$ for large $|\eta|$ in view of the asymptotics (2.2) for P_0 and P_m . Hence, taking the imaginary part of g_2 yields

$$H\psi(\eta) = \begin{cases} x - 4^m(2m+1)t\eta^{2m} & 0 < \eta < \sqrt{u_3} \\ x - 4^m(2m+1)t\eta^{2m} - \frac{[-xP_0(\eta^2, 1, u_2, 0) + 4^m(2m+1)tP_m(\eta^2, 1, u_2, 0)]\eta}{\sqrt{(1-\eta^2)(u_2-\eta^2)(\eta^2-u_3)}} & \sqrt{u_3} < \eta < \sqrt{u_2} \\ x - 4^m(2m+1)t\eta^{2m} & \sqrt{u_2} < \eta < 1 . \end{cases}$$

We then have

$$(4.7) \quad L\psi(\eta) = \begin{cases} x\eta - 4^m t \eta^{2m+1} & 0 < \eta < \sqrt{u_3} \\ x\eta - 4^m t \eta^{2m+1} - \int_{\sqrt{u_3}}^{\eta} \frac{[-xP_0 + 4^m(2m+1)tP_m]\mu}{\sqrt{(1-\mu^2)(u_2-\mu^2)(\mu^2-u_3)}} d\mu & \sqrt{u_3} < \eta < \sqrt{u_2} \\ x\eta - 4^m t \eta^{2m+1} & \sqrt{u_2} < \eta < 1 , \end{cases}$$

where we have used

$$\int_{\sqrt{u_3}}^{\sqrt{u_2}} \frac{[-xP_0 + 4^m(2m+1)tP_m]\mu}{\sqrt{(1-\mu^2)(u_2-\mu^2)(\mu^2-u_3)}} d\mu = 0 ,$$

which is a consequence of (2.3) for P_0 and P_m .

To verify (4.4), we derive an integral formula for ψ . We use (2.28) and (3.1) to calculate

$$\psi_x(\eta) = \begin{cases} \frac{\eta P_0(\eta^2, 1, u_2(x, t), u_3(x, t))}{\sqrt{(1-\eta^2)(u_2(x, t)-\eta^2)(u_3(x, t)-\eta^2)}} & 0 < \eta < \sqrt{u_3(x, t)} \\ 0 & \sqrt{u_3(x, t)} < \eta < \sqrt{u_2(x, t)} \\ \frac{-\eta P_0(\eta^2, 1, u_2(x, t), 0)}{\sqrt{(1-\eta^2)(\eta^2-u_2(x, t))(\eta^2-u_3(x, t))}} & \sqrt{u_2(x, t)} < \eta < 1 \end{cases}$$

Integrating yields

$$(4.8) \quad \psi(\eta) = \begin{cases} -\int_x^{\chi(\eta^2)t} \frac{\eta P_0(\eta^2, 1, u_2(y, t), u_3(y, t))}{\sqrt{(1-\eta^2)(u_2(y, t)-\eta^2)(u_3(y, t)-\eta^2)}} dy & 0 < \eta < \sqrt{u_3(x, t)} \\ 0 & \sqrt{u_3(x, t)} < \eta < \sqrt{u_2(x, t)} \\ \int_x^{\chi(\eta^2)t} \frac{\eta P_0(\eta^2, 1, u_2(y, t), 0)}{\sqrt{(1-\eta^2)(\eta^2-u_2(y, t))(\eta^2-u_3(y, t))}} dy & \sqrt{u_2(x, t)} < \eta < 1 \end{cases}$$

where $\chi(\eta^2)$ is defined in Figure 1. The polynomial P_0 of (2.1) is linear in η^2 and has a zero for $u_3 < \eta^2 < u_2$ because of (2.3). P_0 must be positive for $\eta^2 > u_2$ and negative for $\eta^2 < u_3$. This combined with (4.8) proves $\psi \geq 0$; so (4.4) is verified.

We now continue to verify the variational conditions (4.2) and (4.3). Again, we use (2.28) and (3.1) to calculate

$$\begin{aligned} \frac{\partial}{\partial x} \left(\frac{[-xP_0(\mu^2, 1, u_2(x, t), u_3(x, t)) + 4^m(2m+1)tP_m(\mu^2, 1, u_2(x, t), u_3(x, t))]\mu}{\sqrt{(1-\mu^2)(u_2(x, t)-\mu^2)(\mu^2-u_3^2(x, t))}} \right) \\ = \frac{-\mu P_0(\mu^2, 1, u_2(x, t), u_3(x, t))}{\sqrt{(1-\mu^2)(u_2(x, t)-\mu^2)(\mu^2-u_3^2(y, t))}} \end{aligned}$$

Integrating yields

$$\begin{aligned} \frac{[-xP_0(\mu^2, 1, u_2(x, t), u_3(x, t)) + 4^m(2m+1)tP_m(\mu^2, 1, u_2(x, t), u_3(x, t))]\mu}{\sqrt{(1-\mu^2)(u_2(x, t)-\mu^2)(\mu^2-u_3^2(x, t))}} \\ = -\int_{\chi(\mu^2)t}^x \frac{\mu P_0(\mu^2, 1, u_2(y, t), u_3(y, t))}{\sqrt{(1-\mu^2)(u_2(y, t)-\mu^2)(\mu^2-u_3^2(y, t))}} dy \end{aligned}$$

for $\sqrt{u_3(x, t)} < \mu < \sqrt{u_2(x, t)}$. The single integral in (4.7) can then be written as a double integral.

After interchanging integrals and using (2.3) for P_0 , the double integral is simplified as

$$-\int_{\chi(\eta^2)t}^x \int_{\sqrt{u_3(y, t)}}^{\eta} \frac{\mu P_0(\mu^2, 1, u_2(y, t), u_3(y, t))}{\sqrt{(1-\mu^2)(u_2(y, t)-\mu^2)(\mu^2-u_3^2(y, t))}} d\mu dy.$$

The polynomial P_0 is linear in μ^2 and has a zero for $u_3(y, t) < \mu^2 < u_2(y, t)$. In view of (2.3) for P_0 , we must have

$$\int_{\sqrt{u_3(y, t)}}^{\eta} \frac{\mu P_0(\mu^2, 1, u_2(y, t), u_3(y, t))}{\sqrt{(1 - \mu^2)(u_2(y, t) - \mu^2)(\mu^2 - u_3^2(y, t))}} d\mu < 0$$

for $\sqrt{u_3(y, t)} < \eta < \sqrt{u_2(y, t)}$. Hence, the integral in (4.7) is positive and this verifies the variational conditions (4.2) and (4.3).

We finally consider case (3). By Lemma 3.6, the second equation of (3.2) determines u_2 as an increasing function of x/t in the interval $\beta \leq x/t \leq 4^m$.

We write $\psi = Re(g_3)$ for real η , where

$$g_3 = \sqrt{-1}(x - 4^m(2m+1)t\eta^{2m}) + \frac{\sqrt{-1}[-xP_0(\eta^2, 1, u_2, 0) + 4^m(2m+1)tP_m(\eta^2, 1, u_2, 0)]}{\sqrt{(\eta^2 - 1)(\eta^2 - u_2)}}.$$

The function g_3 is analytic in $Im(\eta) > 0$ and $g_3(\eta) \approx O(1/\eta^2)$ for large $|\eta|$ in view of the asymptotics (2.2) for P_0 and P_m . Hence, taking the imaginary part of g_3 yields

$$H\psi(\eta) = \begin{cases} x - 4^m(2m+1)t\eta^{2m} - \frac{-xP_0(\eta^2, 1, u_2, 0) + 4^m(2m+1)tP_m(\eta^2, 1, u_2, 0)}{\sqrt{(1-\eta^2)(u_2-\eta^2)}} & 0 < \eta < \sqrt{u_2} \\ x - 4^m(2m+1)t\eta^{2m} & \sqrt{u_2} < \eta < 1. \end{cases}$$

We then have

$$L\psi(\eta) = \begin{cases} x\eta - 4^m t \eta^{2m+1} - \int_0^\eta \frac{-xP_0 + 4^m(2m+1)tP_m}{\sqrt{(1-\mu^2)(u_2-\mu^2)}} d\mu & 0 < \eta < \sqrt{u_2} \\ x\eta - 4^m t \eta^{2m+1} & \sqrt{u_2} < \eta < 1, \end{cases}$$

where we have used

$$\int_0^{\sqrt{u_2}} \frac{-xP_0(\mu^2, 1, u_2, 0) + 4^m(2m+1)tP_m(\mu^2, 1, u_2, 0)}{\sqrt{(1-\mu^2)(u_2-\mu^2)}} d\mu = 0,$$

which is a consequence of (2.3) for P_0 and P_m .

The variational conditions (4.2-4.3) and the constraint (4.4) can be verified using the method for case (2).

□

Acknowledgments. We thank Tamara Grava for some of the ideas leading to the proof of Lemma 3.3. V.P. was supported in part by NSF Grant DMS-0135308. F.-R. T. was supported in part by NSF Grant DMS-0404931.

REFERENCES

1. B.A. Dubrovin and S.P. Novikov, “Hydrodynamics of Weakly Deformed Soliton Lattices. Differential Geometry and Hamiltonian Theory”, Russian Math. Surveys 44:6 (1989), 35-124.
2. T. Grava and F.R. Tian, “The Generation, Propagation and Extinction of Multiphases in the KdV Zero Dispersion Limit”, Comm. Pure Appl. Math. 55(2002), 1569-1639.
3. A.V. Gurevich and L.P. Pitaevskii, “Non-stationary Structure of a Collisionless Shock Wave”, Soviet Phys. JETP 38 (1974), 291-297.
4. I.M. Krichever, “The Method of Averaging for Two-dimensional ‘Integrable’ Equations”, Functional Anal. App. 22 (1988), 200-213.
5. P.D. Lax and C.D. Levermore, “The Small Dispersion Limit for the Korteweg-de Vries Equation I, II, and III”, Comm. Pure Appl. Math. 36 (1983), 253-290, 571-593, 809-830.
6. P.D. Lax, C.D. Levermore and S. Venakides, “The Generation and Propagation of Oscillations in Dispersive IVPs and Their Limiting Behavior” in Important Developments in Soliton Theory 1980-1990, T. Fokas and V.E. Zakharov eds., Springer-Verlag, Berlin (1992).
7. P.G. LeFloch, *Hyperbolic Systems of Conservation Laws*, Lectures in Mathematics, Birkhauser, 2002.
8. C.D. Levermore, “The Hyperbolic Nature of the Zero Dispersion KdV Limit”, Comm. P.D.E. 13 (1988), 495-514.
9. V. Pierce and F.R. Tian, “Self-Similar Solutions of the Non-Strictly Hyperbolic Whitham Equations”, Comm. Math. Sci. 4(2006), 799-822.
10. F.R. Tian, “Oscillations of the Zero Dispersion Limit of the Korteweg-de Vries Equation”, Comm. Pure Appl. Math. 46 (1993), 1093-1129.
11. F.R. Tian, “On the Initial Value Problem of the Whitham Averaged System”, in Singular Limits of Dispersive Waves, N. Ercolani, I. Gabitov, D. Levermore and D. Serre eds., NATO ARW series, Series B: Physics Vol. 320, Plenum, New York (1994), 135-141.
12. F.R. Tian, “The Whitham Type Equations and Linear Overdetermined Systems of Euler-Poisson-Darboux Type”, Duke Math. Jour. 74 (1994), 203-221.

13. S.P. Tsarev, “Poisson Brackets and One-dimensional Hamiltonian Systems of Hydrodynamic Type”, Soviet Math. Dokl. 31 (1985), 488-491.
14. S. Venakides, “The Zero Dispersion Limit of the KdV Equation with Nontrivial Reflection Coefficient”, Comm. Pure Appl. Math. 38 (1985), 125-155.
15. G.B. Whitham, “Non-linear Dispersive Waves”, Proc. Royal Soc. London Ser. A 139 (1965), 283-291.

DEPARTMENT OF MATHEMATICS, OHIO STATE UNIVERSITY, 231 W. 18TH AVENUE, COLUMBUS, OH 43210

E-mail address: `vpierce@math.ohio-state.edu`

DEPARTMENT OF MATHEMATICS, OHIO STATE UNIVERSITY, 231 W. 18TH AVENUE, COLUMBUS, OH 43210

E-mail address: `tian@math.ohio-state.edu`